

CHANGING THE POSITION BASIS WITH A UNITARY TRANSFORMATION

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Shankar, R. (1994), *Principles of Quantum Mechanics*, Plenum Press. Section 7.4, Exercise 7.4.9.

The standard representation of the position and momentum operators in the position basis is

$$(1) \quad X \rightarrow x$$
$$(2) \quad P \rightarrow -i\hbar \frac{d}{dx}$$

It turns out it's possible to modify this definition by adding some arbitrary function of position $f(x)$ to P so we have

$$(3) \quad X' \rightarrow x$$
$$(4) \quad P' \rightarrow -i\hbar \frac{d}{dx} + f(x)$$

Since any function of x commutes with X , the commutation relations remain unchanged, so we have

$$(5) \quad [X', P'] = i\hbar$$

Another way of interpreting this change in operators is by using the unitary transformation of the X basis, in the form

$$(6) \quad |x\rangle \rightarrow |\tilde{x}\rangle = e^{ig(X)/\hbar} |x\rangle = e^{ig(x)/\hbar} |x\rangle$$

where

$$(7) \quad g(x) \equiv \int^x f(x') dx'$$

The last equality in 6 comes from the fact that operating on $|x\rangle$ with any function of the X operator (provided the function can be expanded in a power series) results in multiplying $|x\rangle$ by the same function, but with the operator X replaced by the numeric position value.

To verify this works, we can calculate the matrix elements of the old X and P operators in the new basis. We have

$$(8) \quad \langle \tilde{x} | X | \tilde{x}' \rangle = \left\langle x \left| e^{-ig(x)/\hbar} X e^{ig(x')/\hbar} \right| x' \right\rangle$$

At this stage, since the two exponentials are numerical functions and not operators, we can take them outside the bracket to

$$(9) \quad \langle \tilde{x} | X | \tilde{x}' \rangle = e^{-ig(x)/\hbar} e^{ig(x')/\hbar} \langle x | X | x' \rangle$$

$$(10) \quad = e^{-ig(x)/\hbar} e^{ig(x')/\hbar} x' \delta(x - x')$$

$$(11) \quad = x \delta(x - x')$$

The exponentials cancel in the last line since the delta function is non-zero only when $x = x'$.

The above result can also be obtained by inserting a couple of identity operators into 8:

$$(12) \quad \langle x | e^{-ig(x)/\hbar} X e^{ig(x')/\hbar} | x' \rangle = \int \int \langle x | e^{-ig(x)/\hbar} | y \rangle \langle y | X | z \rangle \langle z | e^{ig(x')/\hbar} | x' \rangle dy dz$$

$$(13) \quad = \int \int \langle x | e^{-ig(x)/\hbar} | y \rangle z \delta(y - z) \langle z | e^{ig(x')/\hbar} | x' \rangle dy dz$$

$$(14) \quad = \int \langle x | e^{-ig(x)/\hbar} | z \rangle z \langle z | e^{ig(x')/\hbar} | x' \rangle dz$$

$$(15) \quad = \int e^{i[g(x') - g(x)]/\hbar} \langle x | z \rangle z \langle z | x' \rangle dz$$

$$(16) \quad = \int e^{i[g(x') - g(x)]/\hbar} \delta(x - z) z \delta(z - x') dz$$

$$(17) \quad = e^{i[g(x') - g(x)]/\hbar} x' \delta(x - x')$$

$$(18) \quad = x \delta(x - x')$$

The momentum operator works as follows. Using the original definition 2 on the modified basis we have

$$(19) \quad \langle \tilde{x} | P | \tilde{x}' \rangle = -i\hbar \left\langle x \left| e^{-ig(x)/\hbar} \frac{d}{dx'} e^{ig(x')/\hbar} \right| x' \right\rangle$$

$$(20) \quad = -i\hbar \left\langle x \left| e^{-ig(x)/\hbar} \frac{i}{\hbar} e^{ig(x')/\hbar} \frac{dg(x')}{dx'} \right| x' \right\rangle -$$

$$(21) \quad i\hbar \left\langle x \left| e^{-ig(x)/\hbar} e^{ig(x')/\hbar} \frac{d}{dx'} \right| x' \right\rangle$$

From 7 we have

$$(22) \quad \frac{dg(x)}{dx} = \frac{d}{dx} \int^x f(x') dx' = f(x)$$

This gives

$$(23) \quad \langle \tilde{x} | P | \tilde{x}' \rangle = \left\langle x \left| e^{i[g(x')-g(x)]/\hbar} \left[f(x') - i\hbar \frac{d}{dx'} \right] \right| x' \right\rangle$$

$$(24) \quad = e^{i[g(x')-g(x)]/\hbar} \left[f(x') - i\hbar \frac{d}{dx'} \right] \langle x | x' \rangle$$

$$(25) \quad = e^{i[g(x')-g(x)]/\hbar} \left[f(x') - i\hbar \frac{d}{dx'} \right] \delta(x-x')$$

$$(26) \quad = \left[f(x) - i\hbar \frac{d}{dx} \right] \delta(x-x')$$

This shows that by a unitary change of X basis 6, we transform the position and momentum operators (well, just the momentum operator, really) according to 3. We've multiplied the original $|x\rangle$ states by a phase factor which depends on some function $f(x)$. This doesn't change the matrix elements of X , but it does add $f(x)$ to the matrix elements of P . The commonly used definition of P is thus with $f(x) = 0$.